# Heats of Dilution of Some Aqueous Rare Earth Electrolyte Solutions at 25 $^{\circ}$ C. 1. Rare Earth Perchlorates

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The heats of dilution of aqueous  $La(ClO_4)_3$ ,  $Pr(ClO_4)_3$ , Nd(ClO<sub>4</sub>)<sub>3</sub>, Sm(ClO<sub>4</sub>)<sub>3</sub>, Gd(ClO<sub>4</sub>)<sub>3</sub>, Dy(ClO<sub>4</sub>)<sub>3</sub>, Er(ClO<sub>4</sub>)<sub>3</sub>, and Lu(ClO<sub>4</sub>)<sub>3</sub> solutions have been determined from approximately 0.001 *m* to saturation at 25 °C. The integral heats of solution of La(ClO<sub>4</sub>)<sub>3</sub>·8H<sub>2</sub>O, Nd(ClO<sub>4</sub>)<sub>3</sub>·8H<sub>2</sub>O, Gd(ClO<sub>4</sub>)<sub>3</sub>·8H<sub>2</sub>O, and Er(ClO<sub>4</sub>)<sub>3</sub>·8H<sub>2</sub>O in water at 25 °C have also been measured. The heat of dilution data are represented by empirical polynomial equations, and relative apparent and partial molal heat contents are calculated. The heat content data are correlated with the inner sphere water coordination numbers of the rare earth cations.

This work was undertaken as part of an extensive program to gain a better understanding of rare earth electrolyte solutions through the measurement of their thermodynamic and transport properties (4, 8, 10, 16, 18). Heats of dilution provide a measure of the energy changes involved in complex formation, hydration, modification of the solvent by the hydrated ions, and the electrical work between different solution states. The heats of dilution of a number of rare earth chlorides and nitrates have been measured up to 0.2 m (11, 13). These heats of dilution were shown to conform to the Debye-Hückel limiting law at concentrations below 0.006 m. At higher concentrations the heat data for the chlorides showed a two-series trend across the rare earth series which was correlated with dilute partial molal volume data of the rare earth chlorides (12). The two series trends have been interpreted as reflecting a difference in the inner sphere water coordination between the light and heavy rare earth cations (14). Heats of dilution of aqueous La(ClO<sub>4</sub>)<sub>3</sub> have recently been reported by Vanderzee and Nutter (19). Here we report heat of dilution measurements for eight rare earth perchlorate solutions from 0.001 m to saturation at 25 °C.

# **Experimental Section**

**Apparatus.** The adiabatically jacketed differential calorimeter used in this work has been previously described (11, 13). Although several modifications have been made since then (3, 7), the essential characteristics of the calorimeter remain unchanged. The sample holders previously described (11) were replaced by thin-walled, annealed Pyrex bulbs (2) which were broken by lowering them against pointed stainless steel rods.

To monitor the accuracy of the calorimeter, the heat of neutralization of HCI with NaOH at an ionic strength of 0.003 was determined. From a total of 28 neutralizations an average value of  $-13.338 \pm 0.020$  kcal mol<sup>-1</sup> was obtained at infinite dilution. This is in good agreement with -13.34 kcal mol<sup>-1</sup> recommended by Hepler and Woolley (5).

**Materials.** The solutions were prepared from the rare earth oxides and perchloric acid. The oxides were purified by ion-exchange methods by the Rare Earth Separation Group of the Ames Laboratory. Total rare earth impurities were less than 0.1% by weight, and total other impurities were also less than 0.1% by weight. The stock solutions were adjusted to the equivalence pH to ensure a 1:3 ratio of rare earth ions to perchlorate ions. The stock solutions, secondary stock solutions, and saturated solutions were analyzed by gravimetric sulfate (14) and EDTA (12) methods. The concentrations of the solutions were known to 0.1% in terms of molality. The dilutions were prepared by weight from the stock solutions and conductivity water having a specific conductance of less than  $1 \times 10^{-6}$  mho cm<sup>-1</sup>, all weights being corrected to vacuum. Detailed descriptions of the preparation and analyses of the solutions can be found in ref 3.

Hydrated crystals were grown from saturated solutions of lanthanum, neodymium, and gadolinium perchlorates at 25 °C. The crystals were dried over magnesium perchlorate until the stoichiometric equilibrium of the hydrate was reached. The rare earth content was determined by EDTA titration so that the rare earth content was known to 0.1%. In every case the crystals were the octahydrate to within 0.06 of a water molecule.

Growth of suitable hydrate crystals from the saturated solution of erbium perchlorate was unsatisfactory. For this reason the hydrated crystals obtained for erbium perchlorate were grown from a solution of  $Er(CIO_4)_3$  containing an excess of perchloric acid. Data from the  $Ce(CIO_4)_3$ -HCIO\_4-H<sub>2</sub>O phase system (25) indicated that the pure 8H<sub>2</sub>O hydrate would precipitate out without being contaminated with acid perchlorate, and analysis of the hydrate obtained from the mixture used indicated that the pure octahydrate was obtained. The ratio of  $Er(CIO_4)_3$ -HCIO<sub>4</sub>-H<sub>2</sub>O, in percent by weight, in the mixture was 35.2:24.7:35.2. The crystals grown from this mixture were washed with chloroform and dried under vacuum until they reached the octahydrate composition. The crystals were analyzed as described above.

**Procedure.** One (10 to 20 mL) or two (4 to 10 mL) sample bulbs were suspended in each calorimeter container and enough conductivity water was weighed into each container to give a total liquid content of 900 g.

When a sample of molality  $m_1$ , containing n' moles of rare earth, was diluted into water to give a solution of molality  $m_2$ , a quantity of heat q' was evolved. When the second sample of molality  $m_1$ , containing n'' moles of rare earth, was diluted into the solution of molality  $m_2$  to give a solution of molality  $m_3$ , a quantity of heat q'' was evolved. The relative apparent molal heat content is related to the evolved heat by the expressions

$$q' = n' [\phi_{\rm L}(m_2) - \phi_{\rm L}(m_1)] \tag{1}$$

$$q'' = (n' + n'')\phi_{\rm L}(m_3) - n''\phi_{\rm L}(m_1) - n'\phi_{\rm L}(m_2)$$
(2)

where  $\phi_{\rm L}(m_i)$  is the relative apparent molal heat content at concentration  $m_i$ . In general, two samples could be diluted in one run at high concentrations (q' and q''), while the large size of the sample bulb for dilute solutions allowed only one dilution per run (q').

Similarly, two samples of the crystal hydrates were used to obtain the heats of solution. Dissolving the first sample containing n' moles of rare earth octahydrate gave a solution of molality  $m_2$  with q' calories of heat being evolved. The second sample containing n'' moles of the rare earth hydrate, dissolved into this solution of molality  $m_2$  gave a solution of molality  $m_3$ , evolving q'' calories of heat. The heat of solution to infinite dilution, -L, is related to the heats evolved by

$$q' = n' [\phi_{\mathsf{L}}(m_2) - \overline{L}^{\cdot}] \tag{3}$$

$$q'' = (n' + n'')\phi_{\rm L}(m_3) - n'\phi_{\rm L}(m_2) - n''\overline{L}.$$
 (4)

Since the vapor pressure of water over rare earth perchlorate solutions decreases with concentration, water will evaporate into the empty volume of a sample bulb when a partially filled bulb is broken. The measured heat quantities were corrected for this effect, as well as for the heat of breaking the glass bulbs, and for daily variations in the ratio of the heat capacities of the two calorimeter containers. All heat of dilution and heat of solution measurements were made at  $25.00 \pm 0.02$  °C and the defined thermochemical calorie, 4.1840 absolute J, was used throughout this work.

# **Calculations and Results**

*Heats of Dilution and Heats of Solution.* The integral heats of dilution,  $\Delta H_{i,f}$ , are given by

$$\Delta H_{1,2} = \phi_{\rm L}(m_2) - \phi_{\rm L}(m_1) = q'/n' \tag{5}$$

$$\Delta H_{1,3} = \phi_{\rm L}(m_3) - \phi_{\rm L}(m_1) = (q' + q'')/(n' + n'')$$
 (6)

Similarly, the integral heats of solution,  $\Delta H_{x,f}$ , are given by

$$\Delta H_{\rm x,2} = \phi_{\rm L}(m_2) - \overline{L} = q'/n' \tag{7}$$

$$\Delta H_{x,3} = \phi_{\rm L}(m_3) - L = (q' + q'')/(n' + n'')$$
(8)

The experimental heats of dilution and solution are given in Tables I and II, respectively. The first set of entries for each salt in Table I refers to dilution of the saturated solution. For those groups of dilutions having the same initial concentration,  $m_i$ , the initial concentration is listed only once. For the final concentrations of each dilution or solution,  $m_f$ , an asterisk denotes sample dilution processes corresponding to eq 2 and 6, and 4 and 8, where  $m_f = m_3$ . These samples were diluted into the solution resulting from the dilution of the immediately preceding sample,  $m_f = m_2$ , corresponding to eq 1 and 5, and 3 and 7.

For the heats of dilution the dilution scheme was designed so that the final concentrations,  $m_2$  and  $m_3$ , covered the very dilute concentration range, <0.01 *m*, while the initial concentration,  $m_1$ , systematically covered the range from saturation to 0.01 *m*. The  $\Delta H_{1,2}$  and  $\Delta H_{1,3}$  values were combined to give heats of dilution in the dilute range

$$\Delta H_{3,2} = \Delta H_{1,2} - \Delta H_{1,3} = \phi_L(m_2) - \phi_L(m_3)$$
(9)

The  $\Delta H_{3,2}$  values and the corresponding  $m_i$  and  $m_f$  values are listed immediately following the  $\Delta H_{1,2}$  and  $\Delta H_{1,3}$  data in Table 1.

Initially the heats of dilution were analyzed by the "short chord–long chord" method of Young and co-workers (22–24) and its modification (9, 21). In this method (see ref 11, for example) the  $\Delta H_{3,2}$ , the "short chords", are used to extrapolate to infinite dilution with the aid of the Debye–Hückel limiting slope by means of a  $P_i$  plot,

$$\frac{\partial \phi_{\rm L}}{\partial m^{1/2}} \simeq \overline{P}_{\rm i} = \frac{\Delta H_{3,2}}{m_2^{1/2} - m_3^{1/2}} = A + Bm^{1/2} + Cm + \dots \quad (10)$$

where A is the limiting slope and  $m^{1/2} = 1/2(m_2^{1/2} + m_3^{1/2})$ . Integration of eq 10 gives  $\phi_{\rm L}$  up to 0.01 m, and  $\phi_{\rm L}$  from 0.01 m to saturation can be calculated from eq 5 and 6.

In our attempts to subsequently fit all of the  $\phi_L$  values over the whole concentration range within experimental precision, including the correct limiting behavior, we found that for the rare earth perchlorates this is possible with polynomials of the kind

$$\phi_{\rm L} = \sum_{j=1}^{7} A_j m^{p_j}$$
(11)

where  $A_1$  is the Debye–Hückel limiting slope with  $p_1 = 1/2$ , and the other  $p_j$  are certain multiples of  $\frac{1}{4}$  in the range  $\frac{3}{4}$  to  $\frac{16}{4}$ . This suggested the possibility of fitting the  $\Delta H_{i,f}$  values directly, omitting the more elaborate  $P_i$  analysis.



Figure 1.  $P_i$  plot for Sm(ClO<sub>4</sub>)<sub>3</sub>: plus,  $\Delta H_{3,2}/(m_2^{1/2} - m_3^{1/2})$ ; line,  $(d\phi_L/dm^{1/2})$  from eq 11.

If we assume that a single  $\phi_{L}$  equation can represent all of the concentration range, then from eq 5, 6, and 9 we have

$$\Delta H_{i,f} = \phi_{L}(m_{f}) - \phi_{L}(m_{i})$$
(12)

and from eq 11

$$\Delta H_{i,f} = \sum_{j=1}^{7} A_j (m_f^{p_j} - m_i^{p_j})$$
(13)

The  $A_i$  can be determined by least-squares methods. Then, by setting  $m_i = 0$ , eq 13 reduces to eq 11. The  $\phi_L$  values calculated from this  $\Delta H$  fit agreed with the  $\phi_L$  values obtained by the  $\overline{P}_i$  analysis to within the experimental precision of the data.

The main strength of the  $P_i$  treatment is that the extrapolation to zero concentration depends mostly on the data near where the limiting law should hold, and is not hindered by depending on data at higher concentrations. For this reason it is very important in our present treatment that the equation chosen is capable of fitting the  $\Delta H_{3,2}$  data in the dilute region accurately. In Figure 1 we show  $P_i$  data for Sm(ClO<sub>4</sub>)<sub>3</sub> based on the  $\Delta H_{3,2}$ values, together with  $(d\phi_L/dm^{1/2})$  calculated from our  $\Delta H$  fit based on all the data. It can be seen that the fit is acceptable. The coefficients,  $A_j$ , and powers,  $p_j$ , to be used with eq 11 and 13 are given in Table III. The differences between the experimental  $\Delta H_{i,f}$  and those calculated from eq 13 are given in the fourth column in Table I, and are plotted for Sm(ClO<sub>4</sub>)<sub>3</sub> in Figure 2 vs.  $m_i^{1/2}$ .

The particular powers,  $p_j$ , in eq 13 for each salt were chosen in the following way. First, it was determined that, in general, powers in multiples of  $m^{1/4}$  gave better fits than powers in multiples of m,  $m^{1/2}$ ,  $m^{1/3}$ ,  $m^{1/6}$ ,  $m^{1/8}$ , or  $m^{1/16}$ . In this and the following, the criterion for goodness of fit was the reduced  $\chi^2$ statistic (1). Second, best fits were obtained with  $p_1 = \frac{1}{2}$ , the



**Figure 2.** Comparison between experimental and calculated  $\Delta H$  from eq 13 for Sm(ClO<sub>4</sub>)<sub>3</sub>: solid circles,  $\Delta H_{1,2}$  and  $\Delta H_{1,3}$ ; open circles,  $\Delta H_{3,2}$ .

Journal of Chemical and Engineering Data, Vol. 22, No. 2, 1977 143

Table I. Heats of Dilution of Some Aqueous Rare Earth Perchlorate Solutions at 25 °C

<u>i</u>	10 <sup>4</sup> m <sub>t</sub>	$-\Delta H_{i,f}$ , cal mol <sup>-1</sup>	Exptl – calcd cal mol <sup>-1</sup>	σ, cal mol <sup>-1</sup>	m <sub>i</sub>	10 <sup>4</sup> m <sub>t</sub>	$-\Delta H_{i,f}$ cal_mol <sup>-1</sup>	Expti – calcd cal mol <sup>-1</sup>	σ, cal mol <sup>-1</sup>
	Lan	thanum Perc	blorate		0.001 802	8.970	53.3	6.0	5.1
4,791	8.9700	7089.5	11.4	3.8	0.001 433	7.409	49.8	2.5	3.7
	18.020**	7036.2	5.3	4.7	0.001 459	7.247	53.3	2.1	3.7
	8.9401	7093.3	7.8	3.8	0.002 625	14.311	54.2	4.8	5.4
	17.944*	7037.8	4.1	4.7	0.002 338	10.074	81.4	-5.2	5.4
4.465	7.2469	6347.7	-9.2	2.8	0.005 011	27.458	64.7	5.0	8.2
	14.592*	6294.4	-11.3	3.4	0.004 /17	24.423	73.8	0.9	8.2
	7.4093	6350.9	-14.0	2.8	0.002 457	11.343	68.7 50.0	3.1	3.6
	14.334*	6301.1	-16.4	3.4	0.002 400	15.917	59.9	4.0	3.0
4.093	10.074	5449.5	15.8	4.0	0.003 238	17 783	66.4	-0.8	3.9
	23.377*	5368.1	21.1	4.9	0.003.678	18 602	70.1	21	3.5
	14.311	5407.4	28.7	4.0	0.003 797	18.481	77.9	-1.3	3.5
0.501	26.245*	5353.2	23.9	4.9	0.006 075	29.398	85.5	1.2	3.7
3.581	24.423	4200.7	14.0	0.0 7 X	0.005 892	30.085	78.4	1.8	3.7
	47.109	4166.9	10.0	7.4	0.005 360	27.092	77.0	2.6	3.1
	27.400 50.112*	4243.0	19.1	0.0	0.006 410	31.697	82.3	3.0	2.4
3 255	12 117	3676.4	-3.8	27	0.009 395	39.854	108.1	2.2	2.8
5.205	24 000*	3616.5	-7.8	2.7	0.009 061	43.705	93.0	1.4	2.8
	11.343	3678.9	-0.8	27					
	24 572*	3610.2	-3.9	3.3		Prase	odymium Pe	rchlorate	
2.902	17.783	2959.6	-3.2	2.9	4 695	4 3586	6622.3	3.4	17
2.002	33.432*	2893.2	-2.4	3.5	4.000	8.3484*	6594.0	-60	24
	15.817	2972.5	-5.3	2.9		4.0220	6625.6	4 1	1.6
	32.376*	2903.1	-8.7	3.5		7.8535*	6601.0	-9.1	2.3
2,565	18,481	2356.9	-3.8	2.6	4.255	5.3322	5488.4	0.6	1.8
	37.970*	2279.0	-2.5	3.2		10.604*	5450.8	-5.2	2.5
	18.602	2356.3	-3.9	2.6		5.5215	5483.5	3.5	1.8
	36.784*	2286.2	-6.0	3.2		10.793*	5448.1	-3.8	2.5
2.238	38.341	1747.0	8.3	2.8	3.862	6.6604		_	_
	18.680	1826.1	5.9	3.2		13.047*	4550.8	5.5	2.1
	30.085	1779.1	3.7	2.7		6.1428	4606.5	0.6	1.8
	58.921*	1700.7	1.9	3.4		12.419*	4567.1	-7.0	2.5
	29.398	1782.3	3.1	2.7		7.1657	_	_	_
	60.746*	1696.8	1.9	3.4		14.268*	4544.0	5.3	2.3
1.982	27.092	1432.3	6.4	2.3		7.1976	4595.0	2.5	2.0
	53.597*	1355.3	3.8	2.8		14.370*	4544.3	4.5	2.8
	27.217	1438.0	0.2	2.3	3.398	6.7974	3666.3	-3.4	1.7
1.733	31.697	1123.5	3.2	2.0		15.531*	3612.2	-7.8	2.7
	64.096*	1041.2	0.2	2.1		8.5256			
	31.517					17.094*	3603.8	-7.3	2.2
	63.952*	1040.0	1.7	2.6	2.930	12.332	2751.7	0.8	2.0
1.465	43.705	837.3	-2.6	2.1		23.274*	2700.7	-1.6	2.6
	90.611*	744.3	-3.9	2.6		11.859	2756.5	-1.0	2.0
	39.854	847.0	-1.3	2.1		24.636*	2698.7	-5.0	2.9
4 470	93.954	738.9	-3.5	2.6	2.508	15.362	2022.1	4.8	1.9
1.170	34.281	670.0	-2.5	1.7		31.492*	1956.7	4.5	2.7
1 005	34.200	5710	0.5	1.7		13.005	2016.5	6.9 5.0	1.9
1.005	41.105	571.0	- 1.2	1.7	2 150	19 209	1401 4	5.0	∠.0 1.7
0 816 7	40.322	545.0	-18	1.7	2.150	38 9 17 *	1491.4	- 1.0	1.7
0.0107	32 994	544.0	-0.3	1.3		17 525	1415.5	-3.5	2.4
0 644 8	33 212	5210	0.6	1.3		38 164*	1420.6	-27	24
0.0110	32,160	524.0	1.3	1.3	1.966	25.514	1227 4	-0.4	19
0.493 9	29.441	529.0	1.9	1.6		54.304*	1143.5	2.6	2.7
	26,173	538.0	5.6	1.6		24,422	1230.6	0.6	1.9
0.350 4	26.235	546.0	-2.7	1.8		52.341*	1146.4	4.1	2.6
	23.571	556.0	- 1.5	1.8	1.773	32.347	991.8	-2.3	1.9
0.250 0	16.346	588.0	-4.4	2.7		67.934*	904.9	0.2	2.6
	16.233	586.0	-1.8	2.7		31.656	996.6	-4.8	1.9
0.161 1	11.635	594.0	-3.6	3.2		68.556*	906.6	-2.7	2.7
	11.263	599.0	-5.9	3.2	1.521	44.733	726.8	0.5	1.8
	12.348	585.0	0.5	3.2		86.942*	647.2	-0.9	2.3
0.089 12	7.9919	562.0	4.8	3.7		41.083	729.9	6.9	1.8
	7.9750	560.0	6.9	3.7		87.263*	646.3	-0.5	2.4
0.039 91	3.3636	522.0	-0.7	10.0	1.305	36.255	610.7	-5.2	1.3
0.040.0-	3.1684	526.0	-1.6	10.0		74.613*	522.2	-1.0	1.7
0.010 27	0.8214	378.0	5.7	10.0		36.232	606.4	-0.8	1.3
0.001 794	8.940	55.5	3.7	5.1		74.013 <b>*</b>	526.1	-3.9	1.7

Table I (Continued)

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	104	$-\Delta H_{i,f},$	Expti - calcd	$\sigma$ ,		104	$-\Delta H_{i,f}$	Expti - calcd	$\sigma$ ,
m	10* m <sub>f</sub>				m <sub>i</sub>	10" m <sub>f</sub>	cal mol	calmoi	
1.119	49.791	483.3	-0.8	1.3	2.879	14.130	2563.9	-2.4	2.2
	100.293	396.0	-1.2	1.7		28.580*	2499.6	-1.1	2.7
	51.468	473.7	5.0	1.3		13.242	2573.3	-6.7	2.2
	102.515*	387.9	4.0	1.7		27.889*	2503.9	-3.0	2.7
0.952 9	38.593	461.3	0.0	1.3	2.603	18.992	2074.4	4.2	2.5
	70.345	381.0	-0.3	2.0		41.204*	1996.6	5.2	3.0
	37.905	403.2	-3.3	1.3		19.705	2072.9	2.5	2.5
0 709 6	58 126	380.3	5.5	2.5		40.437*	1998.7	5.2	3.0
0.700 0	58 280	379.3	6.4	1.6	2.254	20.784	1550.5	0.7	1.8
0.504 2	47.716	422.8	-4.8	1.5		40.183*	1480.0	5.2	2.3
	45.822	424.6	-2.0	1.4		20.268	1556.0	-2.5	1.8
	44.984	429.2	-4.5	1.4		39.640	1405.0	0.9	2.3
0.382 0	32.523	474.1	-1.3	1.5	1.964	27.626	1154.1	2.6	2.0
	33.342	468.1	2.1	1.4		58.706	1069.2	4.0	2.4
0.256 4	31.696	484.3	0.8	1.5		29.432	1066.2	3.1	2.0
	31.653	482.4	2.9	1.5	4.070	59.151	1000.2	5.4	2.4
0.160 8	20.866	516.3	-1.1	2.4	1.670	33.339	841.8	-5.3	1.7
	20.794	512.5	3.0	2.4		/1.9/8*	750.3	-3.0	2.1
0.099 76	14.491	517.5	-5.0	3.1		35.117	034.0 749.4	-3.8	1.7
0 0 4 4 7 4	14.664	512.3	0.7	3.1		72.042	740.4	2.2	2.1
0.04174	6.7044	470.1	-6.7	4.0	1.444	52.722	609.9	-1.2	1.7
0 000 935	0.7000	470.2	-0.9	4.0		102.941	524.3	-0.7	2.1
0.000 785	4.339	20.3	3.4 13.2	2.0		107 952*	517.3	-0.2	2.1
0.001.060	5 332	37.7	57	2.4	4 007	107.952	517.3	-0.2	2.1
0.001.079	5.521	35.3	7.4	2.7	1.227	49.323	490.0	4.2	1.7
0.001 242	6.143	39.3	7.7	2.7		47.748	499.0	-1.0	1.7
0.001 437	7.198	50.7	-1.9	3.1	1.005	38.963	445.0	-1.1	1.3
0.001 553	6.797	54.2	4.3	2.8		41.216	437.0	0.7	1.3
0.002 327	12.332	51.0	2.5	2.9	0.794 5	46.717	392.0	1.3	1.4
0.002 464	11.859	57.8	3.9	3.2		50.937	382.0	1.3	1.4
0.003 149	15.362	65.4	0.4	2.9	0.639 0	32.856	431.0	1.8	1.6
0.003 127	15.665	59.6	3.9	2.8		34.928	428.0	-1.6	1.6
0.003 892	18.208	77.9	-4.1	2.6	0.476 3	27.731	473.0	-5.8	2.0
0.003 816	17.525	75.7	-0.7	2.6		26.564	474.0	-2.6	2.0
0.005 430	25.514	83.8	-2.9	2.9	0.360 6	23.387	501.0	-0.4	2.2
0.005 234	24.422	84.2	-3.5	2.8		24.030	495.0	3.1	2.2
0.006 793	32.347	86.9	-2.4	2.9	0.249 6	15.117	549.0	1.7	2.9
0.000 800	31.000	90.1	-2.3	2.9		16.185	542.0	3.1	2.9
0.008 726	44.733	79.0	7.4	2.5	0.159 9	11.149	569.0	-2.1	3.4
0.007 461	36 255	88.5	-4.2	1.8		10.890	568.0	0.7	3.4
0.007 401	36 232	80.3	3.1	1.0	0.100 4	6.4161	563.0	9.3	3.9
0.010.03	49.791	87.3	0.4	1.9		6.7808	562.0	7.0	3.9
0.010 25	51.468	85.7	1,1	1.9	0.042 21	3.6328	517.0	- 16.0	10.0
0.007 634	38.593	80.3	0.2	2.7	0.010.50	3.1082	497.0	11.2	10.0
0.007 602	37.905	78.6	3.3	2.6	0.010 50	0.0947	416.0	- 56.6	10.0
					0 002 023	11 785 <sup>b</sup>	400.0	-47.1	49
	Neo	dymium Pero	blorate		0.001.818	8 602	62.3	-5.3	4.9
4.685	8.6025	6383.0	-1.5	3.6	0.001.746	8 976	48.8	1.9	4.2
	18.182*	6320.7	3.8	4.5	0.001 752	8.726	53.0	-0.1	4.2
	11.785	6343.7	15.6	3.6	0.002 905	14.977	55.4	4.6	5.7
	20.232*	6300.2	14.9	4.5	0.002 982	14.969	65.8	-3.1	5.7
4.509	8.7261	5937.4	-12.0	3.1	0.003 409	17.548	62.9	0.3	5.4
	17.523*	5884.4	-11.9	3.8	0.003 449	17.073	67.5	-0.7	5.4
	8.9760	5932.4	-8.9	3.1	0.003 044	15.896	63.4	-3.6	3.8
	17.464*	5883.6	-10.9	3.8	0.002 789	13.242	69.4	-3.7	2.9
4.075	14.969	4869.2	9.2	4.2	0.002 858	14.130	64.3	-1.3	2.9
	29.823*	4803.4	12.3	5.2	0.004 044	19.705	74.2	-2.7	3.4
	14.977	4864.9	13.4	4.2	0.004 120	18.992	77.8	-1.0	3.4
0.005	29.052*	4809.5	8.8	5.2	0.003 964	20.268	70.2	-3.4	2.5
3.625	17.073	3926.5	7.2	4.0	0.004 018	20.784	70.5	-4.5	2.5
	34.492*	3039.0	7.9	4.9	0.005 974	29.452	81.0	-2.3	2.6
	17.040	3921.3 3959 6	9.9	4.U 1 0	0.005 8/1	27.626	84.9	- 1.9	2.6
3 223	15 206	317/ 3	9.0	4.9 0.0	0.007 204	35,11/	85.4	-1.5	2.3
0.2 60	30 437*	3110.9	-0.3	2.0	0.007 196	33.339 50 700	91.5 0 0 0	-2.3	2.3
	14.062	3193.0	-15.1	3.3	0.010.29	52.722	90.2 85 6	-0.5	2.3
				0.0	0.01020	JE., 22	00.0	0.0	2.0

		$-\Delta H_{if}$	Exptl - calcd	σ.			$-\Delta H_{\rm blue}$	Expt! - calcd	σ.
<b>m</b> i	104 m.	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>	m	10 <sup>4</sup> m <sub>t</sub>	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>
					0.001.010	7.050			1.0
	Sar	marium Perch	lorate		0.001 619	7.659	55.2	-0.4	4.2
4 640	5 8300	6275 6	16.6	2.2	0.001 585	7.678	47.5	5.3	4.0
4.040	11675*	6240.2	6.5	2.2	0.001 936	8.992	52.4	7.2	4.3
	5 6933	6306 4	-12.9	3.1	0.002 016	9.810	47.8	9.4	4.4
	11 560*	6266.0	- 12.0	2.2	0.002 194	10.706	51.5	7.2	3.7
4 004	7 6501	5470.4	- 18.0	3.2	0.002 523	12.445	61.4	-0.7	3.4
4.324	7.0091	5478.4	~-0.5	2.0	0.002 549	12.282	70.9	-8.2	3.4
	16.190*	5423.2	-0.1	3.9	0.003 629	18.103	66.4	0.8	3.4
	7.6784	5475.7	2.0	2.6	0.003 556	17.368	69.4	-0.8	3.3
	15.851	5428.2	-3.3	3.8	0.004 862	24.360	71.6	1.3	3.4
4.006	8.9919	4/48.3	7.9	2.7	0.004 913	24.435	70.9	2.9	3.4
	19.360*	4695.9	0.8	4.1	0.005 454	27.329	74.7	0.6	2.9
	9.8104	4741.4	8.8	2.8	0.005 435	27.060	72.9	2.9	2.9
	20.158*	4693.6	-0.5	4.1	0.005 713	28.296	80.6	-3.2	2.4
3.433	11,142				0.005 730	28.517	78.0	-1.0	2.4
	22.168*	3555.4	-9.3	2.7	0.006 411	31.853	79.7	-0.2	2.6
	10.706	3606.2	-0.5	2.4	0.005 615	28.074	77.1	- 1.0	1.9
	21.936*	3554.8	-7.7	3.4	0.005 584	27.886	76.1	0.0	1.8
3.051	12.445	2913.6	-7.2	2.2		Gad	Iolinium Percl	hlorate	
	25.233*	2852.2	-6.5	3.1	4.611	8.4681	6799.0	2.1	4.0
	12.282	2909.5	-2.1	2.2		8.3175	6807.0	-4.7	4.0
	25.487*	2838.6	6.1	3.2	4.205	6.9064	5852.2	4.0	2.4
2.560	18.103	2072.6	3.0	2.3		13.631*	5802.1	5.7	3.0
	36.291*	2006.2	2.2	3.1		6.8225	5862.0	-5.1	2.8
	17.368	2078.6	0.6	2.2	3.827	7.9580	5035.6	1.0	2.5
	35.564*	2009.2	1.4	3.1		16.176*	4980.8	2.4	3.0
2.244	24.360	1596.7	4.2	2.3		8.0826	5040.6	-5.0	2.5
	48.619*	1525.1	2.9	3.1		16.257*	4982.9	-0.1	3.0
	24.435	1594.2	6.4	2.3	3.499	12.180	4352.0	-6.1	3.6
	49.134*	1523.3	3.4	3.2		12,180	4351.9	-6.0	3.1
1.942	27.329	1221.6	0.9	2.0		23.668*	4289.1	-0.9	3.8
	54.538*	1146.9	0.4	2.6	3.104	17.682	3559.6	3.8	3.5
	27.060	1218.9	4.6	20	•••••	33 5 13	3496.6	4.2	4.3
	54.352*	1146.0	17	2.6		17 497	3557 7	6.7	3.5
1.693	28 296	974.6	-49	1.6		33 443*	3495.4	5.6	4 3
1.000	57 129*	894.0	-17	22	2 891	14 723	3200.0	-5.7	4.1
	28 5 17	964.4	4.5	1.6	2.001	14.608	3197.2	-22	3.5
	57 302*	886 4	5.5	2.0		27 720*	3139.7	-37	4.3
	31 330	956.2	3.5	2.2	2 380	21.120	2324 1	4.0	2.8
	31 953	950.2	0.0	1.8	2.505	41 409*	2024.1	4.0	3.4
	64 109+	937.7	0.0	2.4	2.074	24 621	19/0 3	5.7	27
1 4 4 4	09.074	780.6	-60	1.7	2.074	40.365*	1775 1	37	33
1.444	20.074	700.0	-0.0	1.3		49.303	1953.2	2.1	27
	07 006	703.0	-3.1	1.7		40.351*	1774.0	2.1	2.7
	27.000	770.2	-2.9	1.3	1 000	49.331	1503.4	-67	2.5
1 000	55.636	702.1	-2.9	1.7	1.025	29.724	1445 6	-6.1	2.0
1.200	50.546	015.7	0.2	2.2		00.600	1500.4	-56	2.5
	50.595	615.4	0.4	2.2		29.692	1522.4	-5.6	2.0
1.107	53.342	526.0	3.2	2.0	4 550	57.800	1444.1	-4.2	3.1
	55.428	523.0	1.6	2.1	1.556	14.846	1299.5	-6.5	1.7
0.922 6	74.409	431.6	3.1	2.3		24.275	1258.1	-9.4	2.1
	75.341	431.0	2.1	2.3		14.907	1293.4	-0.7	1.7
0.747 4	75.504	403.3	4.0	2.2		24.433*	1251.9	3.9	2.1
	77.632	398.9	4.8	2.2		38.900	1198.1	0.9	1.7
0.611 6	71.874	406.4	1.4	2.1		74.667*	1115.2	3.0	2.1
	70.637	407.2	2.8	2.1		39.075	1194.7	3.8	1.7
0.498 1	56.991	439.8	-1.1	1.8		74.822*	1113.9	4.1	2.1
0.498 1	57.219	438.5	-0.3	1.8	1.320	42,107	981.9	0.6	2.0
0.375 2	47.588	465.2	0.2	1.6		89.302*	884.2	2.1	2.4
	51.519	457.5	-1.4	1.7		41.977	—		
0.249 7	32.099	508.8	-0.6	1.2		89.246*	884.5	1.9	2.1
	26.838	528.7	-2.5	1.7		21.884	1052.3	-1.0	2.0
	28.641	518.6	1.2	1.6		32.422*	1015.4	-3.8	2.4
0.211 9	28.533	516.5	-0.7	1.6	1.081	20.702	894.8	-1.0	1.3
	30.518	509.1	-0.1	1.4		32.138*	850.6	-0.7	1.6
0.143 6	22.164	521.5	-2.2	2.2		20.967	892.8	-0.2	1.3
	23.648	512.1	1.2	2.1		31.956*	852.9	-2.4	1.6
0.092 45	15.723	509.0	0.7	3.0		43.957	809.0	5.8	1.5
	16.001	501.6	6.7	3.0		43.957	812.0	2.8	1.5
0.048 85	8.4845	475.3	1.1	3.8	0.918 2	56.010	694.0	7.7	2.2
	8.4092	480.4	-3.4	3.8		56.355	696.0	5.0	2.2
0.001 168	5.831 <sup>b</sup>	35.4	10.1	3.4		30.669	767.0	4.1	2.2
0.001 156	5.683	40.5	5.8	3.4		30.980	768.0	2.1	2.2

		$-\Delta H_{i,t}$	Expti - calcd	σ,			$-\Delta H_{i,f}$	Expti - caicd	σ,
m	10 <sup>4</sup> m <sub>f</sub>	cal mol-1	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>	m	104 m.	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>
0.700.0					2.086	10 2 16	2152.3	0.5	2.5
0.733 0	30.492	596.0	5,4	1.5	2.000	38 258+	2086.7	-11	2.5
	30.437	703.0	-1.4	1.5		10 630	2060.7	-0.1	3.4
	25.261	717.0	3.7	1.5		10.030	2130.1	-0.1	2.4
0 550 <b>5</b>	24.800	732.0	-9.5	1.5	1 700	10 400	2000.4	-0.7	3.3
0.552 5	38.465	629.0	0.5	1.5	1.792	19.400	1/28.7	0.8	2.0
	38.490	632.0	-2.6	1.5		39.588	1656.9	2.0	2.8
	24.315	680.0	-2.2	1.5		19.070	1/31.0	0.0	2.0
	24.384	677.0	0.5	1.5		38.574*	1662.6	-0.9	2.8
0.425 6	47.665	580.0	0.4	1.7	1.680	22.096	1572.5	1.7	2.1
	47.790	582.0	- 1.9	1.7		44.459	1499.2	3.1	2.8
	32.558	625.0	-1.3	1.7		21.991	1579.4	-4.8	2.1
	32.627	625.0	-1.6	1.7		43.943*	1506.6	-2.9	2.8
0.327 5	38.143	589.0	-0.9	1.5	1.456	22.059	1317.5	-1.6	1.7
	38.155	589.0	-0.9	1.5		21.831	1321.6	-4.7	1.8
	27.552	621.0	1.9	1.5		45.171*	1242.1	0.0	2.4
	27.426	625.0	-1.6	1.5	1.190	23.448	1061.9	-1.9	1.5
0.198 2	28.228	578.0	4.9	1.3		46.767*	987.7	0.2	2.0
	30.305	580.0	-4.4	1.3		23.619	1060.3	-1.0	1.5
	28.633	582.0	-0.5	1.3		47.036*	989.6	-2.4	2.0
0.113 2	14,784	582.0	1.7	2.8	1.069	26.485	952.7	2.2	1.5
	15,265	570.0	11.0	2.8		52.545*	879.7	1.0	2.0
	15.571	575.0	4.4	2.8		26.601	956.4	-2.0	1.5
0.059 15	8.3290	526.0	10.1	3.7		54.491*	875.9	0.5	2.1
	8 3810	535.0	0.7	3.7	0.923 9	26.664	862.4	-1.8	1.4
	8 667 1	525.0	8.4	37		54.708*	781.1	1.2	1.8
0.001.363	6 906 5	50.1		33		26.269	854.6	7.4	1.4
0.001.626	8.083	57.7	-4 9	33		54.371*	773.8	9.2	1.8
0.001.618	7 958	54.8	-14	3.0	0.745.8	38.423	733.2	0.1	1.6
0.007.010	12 180	62.8	-5.1	3.3		76.227*	647.6	3.6	21
0.002 307	14 609	02.0 57.5	-5.1	4.2		37 128	740.2	-3.1	16
0.002 772	17.000	57.5	1.5	4.8		75 236*	656.6	-37	2.1
0.003 344	17.497	62.3	1.0	4.8	0 588 5	40.862	671.0	-3.0	2.1
0.003 351	17.082	63.0	-0.4	4.8	0.000 0	40.002	674.7	-3.0	2.0
0.004 141	21.013	/4./	-4.1	3.8	0 500 0	40.237	611.6	-4.0	2.0
0.004 935	24.552	79.0	-2.5	3.6	0.509 0	53.171	011.0	4.0	2.3
0.004 936	24.621	74.2	2.0	3.6		33.369	010.4	- 1.3	2.4
0.005 787	29.692	78.3	-1.4	3.4		34.727	005.4	1.8	1.7
0.005 805	29.724	77.8	-0.6	3.4		37.908	656.3	-2.1	1.8
0.007 482	39.075	80.8	-0.3	2.3	0.402 1	41.910	617.5	-1.8	1.9
0.007 467	38.900	82.9	-2.2	2.3		46.014	603.9	1.2	2.0
0.002 443	14.907	41.5	3.1	2.3	0.244 1	31.310	599.5	0.6	1.4
0.002 428	14.846	41.4	3.0	2.3		32.842	596.7	-1.5	1.5
0 003 242	21.884	36.9	2.7	2.7	0.163 2	23.187	586.4	0.1	1.8
0.008 930	42.107	97.7	-1.5	2.7		24.461	577.7	3.7	1.7
0.003 196	20.967	39.9	2.2	1.8	0.104 5	16.709	556.2	1.7	2.7
0.003 214	20.702	44.2	-0.3	1.8		16.799	555.5	2.0	2.7
					0.040 46	6.6273	484.5	-5.8	4.0
	_					6.5234	483.1	-3.5	4.0
	Dys	prosium Perc	hiorate		0.001 048	5.218 <sup>b</sup>	30.1	13.6	3.6
4.602	5.2175	7404.8	14.1	2.4	0.001 128	5.641	38.3	6.4	3.5
	10.478*	7374.7	0.6	3.3	0.001 180	5.929	34.4	10.7	3.3
	5.4583		<u></u>		0.001 209	6.075	37.1	8.4	3.4
	10.659*	7393.0	- 19.0	2.7	0.002 013	10,398	47.8	5.0	3.2
4.317	5.6414	6747.9	-0.2	2.3	0.001 919	9.641	48.6	5.2	3.2
	11.280*	6709.6	-6.5	3.3	0.002 661	13.498	59.5	-0.1	3.6
4.044	5.9287	6111.8	7.8	2.2	0.002 704	13.508	60.1	0.9	3.6
	11.799*	6077.4	-2.9	3.1	0.003 224	16.039	65.3	-0.4	3.6
	6.0749	6114.2	4.0	2.3	0.003 260	16.232	65.4	-0.3	3.7
	12.087 •	6077.1	-4.4	3.2	0.003 826	19.316	65.7	1.5	3.6
3.537	6.3788	4991.2	-3.9	1.9	0.003 751	18.638	67.8	0.4	3.6
	6.2727	4991.7	-3.4	1.9	0.003 959	19.400	71.9	- 1.3	3.1
2.869	10.398	3567.7	7.9	2.2	0.003 857	19.070	68.4	0.9	3.0
	20.128*	3519.9	3.0	3.0	0.004 446	22,096	73.2	-1.3	3.1
	9.6409	3582.9	-2.0	2.1	0.004 394	21.991	72.8	-1.8	3.1
	19.186*	3534.2	-7.1	2.9	0.004 517	21.831	79.5	-4.7	2.6
2.561	13.498	2980.8	-1.7	2.4	0.004 677	23,448	74.2	-2 1	22
	26.606*	2921.4	-1.7	3.3	0.004 704	23.619	70.7	1.4	2.2
	13.508	2975.0	4.1	2.4	0.005 255	26.485	72.9	1.2	22
	27.039*	2914.8	3.3	3.4	0.005 449	26 601	80.5	-24	2.2
2.303	16.039	2513.4	3.1	2.4	0.005 471	26 664	81.3	-30	20
	32.242*	2448.1	3.5	3.4	0.005 437	26 269	80.8	-1 R	2.0
	16.232	2516.3	-0.7	2.5	0.007 623	38.423	85 7	-3.6	2.3
	32.599*	2450.9	-0.4	3.4	0.007 524	37.128	83.6	0.5	2.3
								0.0	2.0

		$-\Delta H_{i,t}$	Exptl – calcd	σ,			$-\Delta H_{i,f}$	Exptl - calcd	σ,
mi	10 <sup>4</sup> m <sub>t</sub>	cal mol-1	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>	mi	104 m <sub>t</sub>	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>	cal mol <sup>-1</sup>
	E	rbium Perchk	orate		0.360 3	40.145	603.0	-4.6	1.7
4.627	8.8685	7399.3	0.1	3.0		40.094	599.0	-0.4	1.7
	17.514*	7338.0	10.2	3.6		39.980	599.0	-0.1	1.7
	8.3348	7400.0	3.5	3.0		39.930	600.0	-1.0	1.7
	17.156*	7350.0	-0.1	3.6		33.143	622.0	-3.1	1.7
4.215	15.817	6396.4	-4.7	4.8		33.109	620.0	-1.0	1.7
	32.274*	6323.8	2.3	5.9	0.230 7	6.4567	712.7	-0.9	2.7
	10.824	6431.7	-11.0	4.8		12.496*	666.9	0.7	3.3
	21.716*	6372.1	-7.7	5.9		6.2500	/1/.1	-3.4	2.7
	10.883	6435.2	- 14.9	4.0		12.510	615.4	2.1	3.3
	23,098	6307.8		5.5		22.753	615.0	2.5	3.1
3 413	8 6966	4648.9	-21.0	2.6	0 202 0	25.020	588.0	0.4	16
5.415	18 207*	4579.9	10.2	3.1	0.202 9	19 598	621.0	-2.6	1.6
	8 6554	4651.2	-5.0	2.6		28.966	581.0	0.8	1.6
	18.284*	4575.6	14.1	3.1	0 115 4	14 646	560.0	16.0	2.9
3.086	9.3086	3963.2	-4.5	2.9	0,110 1	15.801	552.0	17.9	2.9
	20.458*	3880.7	16.6	3.5	0.001 716	8.335 <sup>b</sup>	50.0	3.6	4.0
	9,1930	3956.2	3.4	2.9	0.001 751	8.868	61.3	- 10.1	4.0
	19.272*	3890.6	11.9	3.5	0.002 310	10.883	67.4	-5.8	6.5
	9.2294	3965.0	-5.7	3.3	0,002 172	10.824	59.6	-3.3	6.5
	14.493	3929.4	-3.2	2.9	0.003 227	15.817	72.6	-7.0	6.5
	28.794*	3865.8	-0.7	3.5	0.001 828	8.655	75.6	- 19.1	3.5
	14.432	3925.3	1.3	2.9	0.001 821	8.697	69.0	<del></del> 13.2	3.5
	29.074*	3867.6	-3.5	3.5	0.002 907	14.432	57.7	4.7	3.9
2.676	16.540	3113.7	0.9	3.1	0.002 879	14.493	63.6	-2.5	3.9
	33.397*	3043.8	5.4	3.8	0.001 927	9.193	65.6	-8.6	3.9
	16.557	3117.7	-3.1	3,1	0.002 046	9.309	82.5	-21.0	3.9
	33.281*	3052.0	-2.5	3.8	0.003 328	16.557	65.7	-0.7	4.2
2.358	36.518	2468.5	3.8	4.0	0.003 340	16.540	69.9	-4.4	4.2
	08.140	2392.9	0.9	5.9	0.006 810	30.724	75.6	-0.0	6.6
	50.724 68.096*	2409.0	8.1	59	0.006 8 15	30.5 10 40 196	75.0	-3.2	6.0
2 033	39 930	1957.2	14 4	44	0.007 900	39 930	82.5	-0.9	6.0
2.000	78 925*	1874 7	-13.5	5.4	0.007 090	32 753	89.6	-7.3	3.9
	40.196	1959.4	- 17.3	4.4	0.006.719	32,501	82.4	0.5	3.9
	78.997*	1875.2	-14.1	5.4	0.009 197	44.316	95.3	-5.0	4.3
1.700	32,501	1508.3	-0.7	2.9	0.009 264	44,409	94.0	-3.0	4.3
	67,191*	1425.9	- 1.2	3.5	0.011 41	53.978	101.0	-3.9	4.0
	32.753	1508.3	- 1.5	2.9	0.011 54	53.802	102.3	-3.1	4.0
	67.273*	1418.7	5.9	3.5	0.001 251	6.250	51.7	-5.5	3.6
1.498	44.409	1243.5	-4.0	3.2	0.001 250	6.457	45.8	-1.6	3.6
	92.641*	1149.5	-1.0	3.9		Li	utetium Perch	lorate	
	44.316	1244.3	-4.5	3.2	4.634	11.096	7233.1	-4.4	4.7
	91.968*	1149.0	0.5	3.9		21.874*	7175.9	-2.6	5.8
1.227	53.978					11.062	7229.0	-0.1	4.7
	114.426*	867.7	-6.6	3.3	4 000	22.212*	7166.4	5.5	5.8
	53.802	956.7	2.3	3.0	4.039	14.3 19	5851.2	-0.7	5.1
	115.434°	004.4	-29	3.7		29.225	5772.5	12.3	5.1
	114 105*	901.0 860.6	2.9	3.7		30 714*	5759.0	16.8	6.3
	53 451	960.0	-0.2	3.4	3,567	12.076	4829.5	- 12.3	3.7
1.055	59 0 13	815.0	4.3	3.6	0.001	26.657*	4753.1	-3.7	4.5
	56.912	820.0	3.7	3.6		12.795	4815.9	-3.1	3.7
	57.093	825.0	-1.7	3.6		24.621*	4760.6	-3.5	4.5
	57.851	819.0	2.7	3.6	3.196	12.681	4045.3	<del>-</del> 11.0	3.1
	76.668	781.0	5.5	3.6		24.493*	3985.1	-6.7	3.8
	76.825	782.0	4.3	3.6		14.799	4013.5	8.8	3.1
0.840 8	80.407	653.0	6.1	3.2		27.479*	3963.4	3.9	3.8
	79.959	656.0	3.8	3.2	2.870	11.202	3392.0	4.7	2.8
	59.413	697.0	0.2	3.2		12.027	3386.9	4.5	2.4
	59.614	692.0	4.8	3.2	0 5 0 0	24.671*	3327.3	3.5	3.0
	59.506	696.0	1.1	3.2	2.529	20.867	27 13.7	4.3	3.0
0.662 6	35.402	001.U 697.0	5.3 _0.9	2.1		18 250	2000.9	5.0	J.Z
	30.420	007.0	-0.0 	2.1		36 132*	2663.0	1.2	3.9
	61.450	618.0	5.3	2.1	2,232	16.941	2236.3	2.0	2.5
	61.450	620.0	3.3	2.1		31.764*	2178.9	1.0	3.0
0.512 9	32.879	651.0	2.2	1.6	1.980	17.181	1863.7	-6.1	1.6
	32.696	656.0	-2.2	1.6		35.713*	1791.3	-3.1	1.9
	32.684	655.0	-1.2	1.6		17.741	-	_	-
	32.971	656.0	-3.1	1.6		35.617*	1790.7	-2.3	1.7

m <sub>i</sub>	10 <sup>4</sup> m <sub>t</sub>	$-\Delta H_{i,f}$ , cal mol <sup>-1</sup>	Exptl — calcd cal mol <sup>-1</sup>	σ, cal mol <sup>-1</sup>	m <sub>i</sub>	104 <i>m</i> f	–∠ <i>H</i> i,f, cal mol <sup>−1</sup>	Exptl — calcd cal mol <sup>-1</sup>	σ, cal mol <sup>-1</sup>
1.676	23,561	1431.6	0.7	2.0	0.154 6	9.3269	635.0	-0.6	3.5
	47,458*	1358.3	1.2	2.4		8.8209	633.0	5.1	3.5
	22.307	1436.3	1.1	2.0	0.098 80	5.7696	603.0	8.0	3.9
	46.690*	1359.0	2.4	2.4		6.3706	601.0	4.2	3.9
1.432	27.741	1152.8	-1.0	1.9	0.067 17	5.7360	560.0	-2.5	10.0
	55.249*	1074.4	2.4	2.3		4.9640	570.0	-4.6	10.0
	26,988	1153.7	0.8	1.9	0.009 869	0.9187	363.0	-20.4	10.0
	56.355*	1069.9	4.6	2.3	0.002 221	11.062	62.6	-5.6	6.4
1 174	35 653	902.6	1.0	17	0.002 187	11.096 <i><sup>b</sup></i>	57.2	- 1.9	6.4
1.174	64 600*	902.0	3.5	21	0.003 071	14.992	69.6	-4.6	7.0
	32 024	014 3	0.5	2.1	0.002 922	14.319	78.9	-15.2	7.0
	52.024	514.5	0.5	2.0	0.002 462	12.795	55.3	0.5	5.0
1.006	36.216	792.5	-2.1	1.6	0.002 666	12.076	76.4	-8.7	5.0
	69.789*	713.1	0.8	1.9	0.002 748	14.799	50.1	4.9	4.2
	36.192	792.4	-2.0	1.6	0.002 449	12.681	60.2	-4.2	4.2
	71.944*	709.1	1.0	1.9	0.002 467	12.027	59.6	1.0	3.3
0.806 8	54.716	645.0	0.9	2.5	0.003 839	20.867	59.8	0.5	3.6
	53.832	648.0	-0.2	2.5	0.003 176	16.941	57.4	1.0	3.3
0.640 9	51.251	595.0	3.9	2.2	0.003 571	17.181	72.4	-2.9	2.1
	50.623	599.0	1.3	2.2	0.004 669	22.307	77.3	-1.3	2.7
0.494 4	36.084	609.0	-2.8	1.6	0.004 746	23.561	73.3	-0.5	2.7
	37.259	607.0	-4.2	1.6	0.005 636	26.988	83.8	-3.7	2.5
0.358 6	26.615	614.0	0.5	1.3	0.005 525	27.741	78.4	-3.4	2.5
	27.783	610.0	0.3	1.3	0.006 461	35.653	71.0	-2.6	2.3
0.249 6	11.683	656.0	4.9	2.9	0.007 194	36.192	83.3	-2.9	2.1
	13.690	650.0	- 1.1	2.9	0.006 979	36.216	79.4	-3.0	2.1

<sup>a</sup> For a starred sample f = 3 and its corresponding f = 2 value (unstarred) is given immediately above. <sup>b</sup> For each salt, all entries above this point are  $\Delta H_{1,2}$  or  $\Delta H_{1,3}$  values, the rest are  $\Delta H_{3,2}$  values.

Table II. Heats of Solution	of Some Rare	Earth Perchlorate	Hydrates at 25	°C
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Hydrate	10 <sup>4</sup> m <sub>f</sub>	$-\Delta H_{x,b}$ cal mol <sup>-1</sup>	$\phi_{L}(m_{\rm f})$ cal mol <sup>-1</sup>	Lَ <sup>.</sup> , cal mol
La(CIO <sub>4</sub> ) <sub>3</sub> •8H <sub>2</sub> O	5.693	9 305.4	145.2	9450.6
	12.454* <sup>a</sup>	9 239.3	203.0	9442.3
	5.741	9 258.5	145.7	9404.2
	11.861*	9 207.9	198.9	9406.8
			Average	9426 ± 21
Nd(CIO <sub>4</sub> ) <sub>3</sub> •8H <sub>2</sub> O	6.061	9 304.7	134.7	9439.4
_	15.155*	9 232.5	198.3	9430.8
	8.952	9 239.6	159.2	9398.8
	14.123*	9 201.1	192.6	9393.7
			Average	9416 ± 20
Gd(CIO <sub>4</sub> ) <sub>3</sub> •8H <sub>2</sub> O	8.439	10 068.4	157.0	10225.4
	11.951*	10 054.8	182.0	10236.8
	6.381	10 135.9	139.2	10275.1
	13.727*	10 067 2	192.9	10260.1
			Average	10249 ± 19
Er(ClO <sub>4</sub> ) <sub>3</sub> -8H <sub>2</sub> O	3.667	13 333.3	107.4	13440.7
	2.088	13 306.8	83.6	13390.4
			Average	$13416 \pm 26$

<sup>a</sup> For a starred sample f = 3 and its corresponding f = 2 value (unstarred) is given immediately above.

limiting law power in molality. Third, approximately equally good fits were obtained whether  $A_1$  was fixed at the limiting slope, 6990 (3), or allowed to vary. Therefore  $A_1$  was fixed at the limiting law value. Fourth, all the possible fits with combinations of six multiples of  $m^{1/4}$  between  $m^{3/4}$  and  $m^{16/4}$  were performed. The final powers were chosen from those fits with low reduced  $\chi^2$  values, subject to the constraint that adjacent rare earth perchlorates should have similar sets of powers as far as possible.

It should be pointed out that the coefficients in eq 13 cannot be determined using standard least-squares programs, since there are two independent variables,  $m_i$  and  $m_i$ ; specifically, the independent variable terms are not  $(\Delta m)^{p_i}$ . However, linear least-squares regression analysis can still be used if the two independent variables are incorporated into the matrix building steps.

The advantages of the present method of data treatment are twofold. First, the two-step  $\overline{P_i}$  analysis is reduced to a simple and direct one-step fit of the measured  $\Delta H$  values. Second, the whole data set is fitted to one least-squares equation, avoiding discontinuities in  $\phi_{\rm L}$  and derived quantities which occur when the data set is fitted in segments. These advantages depend on the existence of empirical equations that are capable of fitting the data within the statistical errors in the data, in addition to extrapolating to zero concentration correctly. Even if such equations cannot be found, the  $\overline{P_i}$  fit, eq 10, can still be replaced

Salt <sup>a</sup>	Р2 Аз	<b>Р</b> 3 Аз	Р4 Ал	$\rho_5$	<i>р</i> 6 Ас	р <sub>7</sub> Ал
	1.00	1.25	1.50	1.75	2.75	4.00
La(CIO <sub>4</sub> ) <sub>3</sub>	-53 930.420	119 602.8009	-110 785.3187	39 847.259 9	-849.299 91	17.608 558
	0.75	1.25	2.00	3.25	3.75	4.00
Pr(ClO <sub>4</sub> ) <sub>3</sub>	-9 866.816	4 762.0196	-1 736.5413	1 487.968 31	-1 465.468 192	587.561 198
	0.75	1.25	2.00	3.25	3.75	4.00
Nd(CIO <sub>4</sub> ) <sub>3</sub>	-9 789.311	4 543.8264	-1 569.9242	1 354.352 59	-1 330.857 723	532.680 189
	0.75	1.25	2.00	3.25	3.75	4.00
Sm(ClO <sub>4</sub> ) <sub>3</sub>	-9 804.321	4 781.3713	-1 733.9837	1 417.622 61	-1 387.953 071	555.814 673
	0.75	1.25	2.00	3.25	3.75	4.00
Gd(CIO <sub>4</sub> ) <sub>3</sub>	-9 422.495	4 341.3987	-1 245.6835	1 005.995 37	-985.614 470	394.371 910
	0.75	1.25	2.00	2.25	2.50	2.75
Dy(CIO <sub>4</sub> ) <sub>3</sub>	-9 989.583	7 097.0994	-20 609.1495	32 196.835 6	-18 034,138 58	3504.302 75
	0.75	1.25	2.00	2.25	2.50	2.75
Er(ClO <sub>4</sub> ) <sub>3</sub>	-10 126.464	7 609.3830	-23 931.1872	37 654.248 2	-21 236.443 33	4157.844 04
	0.75	1.25	2.00	2.25	2.50	2.75
Lu(CIO <sub>4</sub> ) <sub>3</sub>	-10 064.327	7 102.3011	-20 755.8581	32 350.593 8	-18 042.102 01	3484.477 10

 ${}^{a}\rho_{1} = 0.50, A_{1} = 6990.00$  for all saits.



Figure 3. Relative apparent molal heat content of some aqueous rare earth perchlorates: from eq 11; DHLL, Debye-Hückel limiting law.

by the  $\Delta H$  fit, eq 13, where the  $\Delta H_{i,t}$  values are restricted to the "short chords",  $\Delta H_{3,2}$ . This avoids the error introduced by the approximation to the derivative, indicated in eq 10.

The relative partial molal heat contents of the solute,  $L_2$ , and the solvent,  $\overline{L}_1$ , were calculated from

$$\overline{L}_{2} = \phi_{L} + m \left(\frac{\partial \phi_{L}}{\partial m}\right)_{\mathrm{T},\mathrm{P},\mathrm{n}_{1}}$$
(14)

$$\overline{L}_{1} = \frac{-M_{1}m^{2}}{1000} \left(\frac{\partial\phi_{L}}{\partial m}\right)_{T,P,n_{1}}$$
(15)

where  $M_1$  is the molecular weight of water, 18.0154 g mol<sup>-1</sup>. Graphs of  $\phi_L$ ,  $L_2$ , and  $L_1$  for representative rare earths are shown in Figures 3, 4, and 5.

The heats of solution to infinite dilution fo the rare earth perchlorate octahydrates, -L, were calculated from eq 7 and 8, using  $\Delta H_{x,f}$  from Table II and with  $\phi_{L}(m_2)$  and  $\phi_{L}(m_3)$  calculated from eq 11. The L are listed in Table II.



Figure 4. Relative partial molal heat content of some aqueous rare earth perchlorates: from eq 11; DHLL, Debye-Hückel limiting law.

**Errors.** Random errors in the heat measurements were found to be correlated with the actual heat evolved during a run, *q*. The standard deviation expected for each  $\Delta H$  measurement,  $\sigma$ , based on the heat evolved, *q*, is listed in the fifth column in Table I. These standard deviations in  $\Delta H$  ranged from  $\sigma = 10$  cal mol<sup>-1</sup> for the very few runs where *q* was less than 0.3 cal, to an optimum of about  $\sigma = 1$  cal mol<sup>-1</sup> at q = 1.5 cal, and rising to  $\sigma = 4.5$  cal mol<sup>-1</sup> at q = 6 cal. These standard deviations were used to obtain weighting factors,  $1/\sigma^2$ , for the  $\Delta H$  fits. The standard deviation of the least-squares fits of  $\Delta H$  ranged from 4 to 6 cal mol<sup>-1</sup>. These standard deviations of the fits compare favorably with the standard deviations of the experimental  $\Delta H$  values, and indicate that the systematic deviations of the experimental  $\Delta H$  values from the equations were very small.

In order to determine the propagation of the random errors,  $\sigma$ , in  $\Delta H$ , to the derived quantities  $\phi_{L}$ ,  $\overline{L}_1$ , and  $\overline{L}_2$ , a random variable technique was used (3). In Table IV we list  $\phi_{L}$ ,  $(d\phi_{L}/d\phi_{L})$ 

Table 14. Then touthannic Frudenies and citurs for Latervals Solutions at 20
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Molality	$\phi_{ m L}$ , cal mol $^{-1}$	$\mathrm{d}\phi_{\mathrm{L}}/\mathrm{d}m^{1/2}$	$-\tilde{L}_1$ , cal mol <sup>-1</sup>	$\overline{L}_2$ , cal mol <sup>-1</sup>
0.001 794	235.7 ± 5.1	4 483.2 ± 39.9	0.003 07 ± 0.000 03	330.6 ± 5.2
0.002 625	273.9 ± 5.4	4 141.8 ± 33.4	0.005 02 ± 0.000 04	$380.0 \pm 5.5$
0.003 678	311.4 ± 3.5	3 818.2 ± 27.5	0.007 67 ± 0.000 06	427.2 ± 3.6
0.005 011	$348.5 \pm 8.2$	3 506.7 ± 22.3	0.011 20 ± 0.000 07	472.6 ± 8.2
0.006 410	379.8 ± 2.4	3 250.6 ± 18.2	0.015 03 ± 0.000 08	$509.9 \pm 2.5$
0.009 061	426.1 ± 2.8	2 882.5 ± 13.1	0.022 39 ± 0.000 10	$563.3 \pm 2.9$
0.010 27	443.4 ± 10.0	2 748.0 ± 11.5	0.025 76 ± 0.000 11	582.7 ± 10.0
0.039 91	636.2 ± 10.0	1 361.1 ± 9.0	0.097 75 ± 0.000 64	772.1 ± 10.0
0.089 12	735.0 ± 3.7	712.7 ± 7.6	0.170 8 ± 0.001 8	841.4 ± 3.9
0.161 1	787.7 ± 3.2	345.9 ± 4.9	0.201 5 ± 0.002 8	857.1 ± 3.4
0.250 0	810.6 ± 2.7	$134.4 \pm 4.6$	0.151 4 ± 0.005 2	844.2 ± 2.9
0.350 4	817.2 ± 1.8	22.8 ± 4.9	0.042 6 ± 0.009 2	823.9 ± 2.3
0.493 9	817.2 ± 1.6	$0.1 \pm 4.5$	$0.000 \pm 0.014$	817.3 ± 2.2
0.644 8	821.3 ± 1.3	103.4 ± 3.7	0.482 ± 0.017	862.9 ± 2.0
0.816 7	842.7 ± 1.3	344.9 ± 3.1	$2.293 \pm 0.020$	998.5 ± 1.9
1.005	894.4 ± 1.7	726.8 ± 2.8	6.596 ± 0.025	1 258.7 ± 2.2
1.176	970.8 ± 1.7	1 153.3 ± 2.6	13.248 ± 0.029	1 596.1 ± 2.2
1.465	$1166.5 \pm 2.3$	1 988.2 ± 2.4	31.756 ± 0.039	2 369.7 ± 2.8
1.733	$1421.1 \pm 2.3$	$2831.3 \pm 2.8$	58.184 ± 0.057	3 284.7 ± 2.9
1.982	$1716.0 \pm 2.4$	$3630.5 \pm 3.2$	91.252 ± 0.081	$4251.6 \pm 3.3$
2.238	2071.5 ± 3.0	$4437.6\pm3.4$	133.83 ± 0.10	$5390.8\pm4.0$
2.565	2591.6 ± 2.9	$5415 \pm 3.3$	$200.39 \pm 0.12$	$6928.1 \pm 3.9$
2.902	3191.1 ± 3.2	6 338.7 ± 3.8	282.27 ± 0.17	8 590.2 ± 4.6
3.255	3873.3 ± 3.0	7 210.6 ± 5.4	381.43 ± 0.29	10 377.8 ± 5.7
3.581	4541.7 ± 6.7	7 944.0 ± 6.4	484.91 ± 0.39	$12058.1\pm9.1$
4.093	$5651.0 \pm 4.4$	9 032.3 ± 6.2	$673.72 \pm 0.46$	14 787.8 ± 7.7
4.465	6499.8 ± 3.1	9 861.9 ± 15.6	838.1 ± 1.3	16 919.1 ± 16.8
4.791	$7277.6 \pm 4.2$	10 693.6 ± 35.8	$1010.1 \pm 3.4$	18 980.9 + 39.4





Sm Gd Dy Er

Figure 5. Relative partial molal heat content of water in some aqueous rare earth perchlorates, from eq 11.

Figure 6. Relative apparent molal heat content of some aqueous rare earth perchlorates, from eq 11.

 $dm^{1/2}$ ),  $\overline{L}_1$ , and  $\overline{L}_2$  for La(ClO<sub>4</sub>)<sub>3</sub> as calculated fom eq 11, 14, and 15, together with the standard deviations expected in these quantities. The random errors in  $\phi_L$  are equal to the random errors in  $\Delta H$  for the same concentration. Even though the errors in  $\phi_L$  are rather uniform over the whole concentration range, the errors for the derivative of  $\phi_L$  increase at both ends of the data set. This reflects the fact that the equation is not constrained as well at the endpoints as in the center of the data set.

In addition to the above random errors, the following systematic errors should be noted. The possible error in the calorimeter calibration is about 0.2% in  $\Delta H$  (and  $\phi_L$ ), which amounts to at most 20 cal. The error due to 0.1% error in the concentration of the solutions ranges from being negligible at low

concentrations, to 10 cal at the highest concentrations. From a comparison of the results from the  $P_i$  analysis and the results of various  $\Delta H$  fits tried, we feel the extrapolation to infinite dilution results in a possible error of at most 10 cal in  $\phi_L$ . In summary, the errors in  $\phi_L$ ,  $\overline{L}_2$ , and  $\overline{L}_1$  should be as indicated in Table IV for La(ClO<sub>4</sub>)<sub>3</sub> with the addition of the above possible systematic errors.

The  $\phi_L$  values for La(ClO<sub>4</sub>)<sub>3</sub> reported by Vanderzee and Nutter (*19*) are systematically larger than those reported here, the difference increasing to approximately 18 cal. This difference is well within the estimated errors of either data set.



Figure 7. Relative partial molal heat content of some aqueous rare earth perchlorates, from eq 11.

# Discussion

The general shape of the  $\phi_{L}$  and  $\overline{L}_{2}$  curves is typical of  $\phi_{L}$  and  $\overline{L}_2$  curves for "normal salts" such as the alkali, alkaline earth, and transition element perchlorates (20). As was the case for the dilute solution studies of the rare earth chloride and nitrate solutions (11, 13), the rare earth perchlorates approach the Debye-Hückel theoretical behavior in dilute solutions. This was also found by Vanderzee and Nutter for La(CIO<sub>4</sub>)<sub>3</sub> (19). Between approximately 0.1 and 1 m the  $\phi_1$  curves go through a pronounced inflection point which is also evident for other perchlorate solutions, although the inflection points occur at higher concentrations for the monovalent and divalent perchlorates. Above 1 *m* the  $\phi_1$  curves rise rapidly to very high values, typical of strongly hydrated ions. When the rare earth perchlorates,  $Ca(CIO_4)_2$  (20) and  $NaCIO_4$  (20), are plotted on the same graph, the rare earth salts fall about where one would expect them with respect to charge and size of the rare earth ions as a group.

 $\phi_L$ ,  $L_2$ , and  $L_1$  are plotted at even molalities across the rare earth series in Figures 6, 7, and 8, respectively.  $\phi_1$  and  $L_2$  decrease from La to about Nd, then increase to around Gd-Dy and then decrease again for the rest of the series. This behavior has also been observed in the heat properties of the rare earth chlorides in dilute solutions (11, 13), the partial molal volume properties of the rare earth chlorides (15), perchlorates (17), and nitrates (16), and the expansibilities of the rare earth chlorides (4). The displacement of these properties in the middle of the rare earth series has been attributed to changes in the inner and outer hydration spheres, caused by a decrease in the number of waters in the inner hydration sphere of the rare earth cations (14). Owing to the decreasing size of the rare earth ion, the inner sphere water coordination of Tb to Lu is thought to be lower than the inner sphere water coordination of La to Nd by one water molecule. The ions between Nd and Tb have mixtures of the higher and lower water coordination. The position of the hydration change in the rare earth series at infinite dilution or in very dilute solutions has invariably been between Nd and Tb. At the lowest concentration shown for  $\phi_L$ ,  $L_2$ , and  $L_1$  the position of the break is near these two ions for the rare earth perchlorates.

The fact that the break in the series remains evident over most of the concentration range indicates that the rare earth ions do not form inner sphere rare earth perchlorate complexes, in agreement with commonly held views. However, it does not of



Figure 8. Relative partial molal heat content of water in some aqueous rare earth perchlorates, from eg 11.

course rule out outer sphere ion pairing. The leveling out of  $L_2$ at 4.6 m across the series does appear to be real, since this behavior also occurs in the partial molal volumes of the rare earth perchlorates (17). This leveling out at 4.6 m in  $L_2$  does not appear in  $\phi_{L}$  and points out that the partial molal property often gives information not present in the apparent molal property, particularly at high concentrations.

From graphs of  $\phi_{\rm L}$  for the alkali and alkaline earth halides and perchlorates (20) it becomes immediately evident that  $\phi_{\rm L}$  decreases guite regularly with increasing ion size for a common cation or anion series.  $\phi_{\rm L}$  also increases with cation charge for a given cation size. This relationship holds over the whole concentration range for all the salts, and has been quantified for dilute solutions in the case of monovalent and divalent halides by Leung and Millero (6). The rare earth perchlorate values fall above the divalent and monovalent perchlorates of similar cation size (i.e., Ca(ClO<sub>4</sub>)<sub>2</sub> and NaClO<sub>4</sub>), as expected. However, from Figure 6 it is clear that the  $\phi_{L}$ -cation size relationship within the rare earth series does not hold for the rare earth perchlorates. since, for a given inner sphere hydration series,  $\phi_{\rm L}$  decreases with decreasing cation size. Whether this anomaly holds for all trivalent salt solutions will have to await  $\phi_{L}$  measurements for smaller and larger trivalent cations than the rare earths.

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# Vapor-Liquid Equilibrium Constants of Alkylbenzenes in n-Alkane Solvents at Infinite Dilution by Gas-Liquid Chromatography

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Vapor-liquid equilibrium constants at infinite dilution of benzene and alkylbenzenes in n-octacosane, ndotriacontane, and n-hexatriacontane solvents were measured by gas-liquid chromatography at five temperatures, 78.0, 84.0, 90.0, 96.0, and 102.0 °C. The thermodynamic properties of solution of alkylbenzenes from the gas phase to n-alkane solvent were also determined from the data of vapor-liquid equilibrium constants. The partial molar enthalpy of solution is related with the structure of alkylbenzenes through an equation by a set of values of structural parameters.

The properties of molecules in liquid phases have been greatly studied by experimental methods of static and dynamic techniques. Theoretical approaches of the behavior of molecules in pure liquids and liquid mixtures are mainly limited to ideal or simple liquid systems. Study for theoretical research demands therefore determining the thermodynamic properties of liquid in simple systems. In liquids consisting of pure substance, molecules frequently have association by hydrogen bonding or other strong intermolecular interaction. The molecular association distorts the liquid phase from the ideal to nonideal liquid system, e.g., polyassociation of alcohols. In solutions at infinite dilution, solute molecules no longer associate with each other, but associate with surrounding solvent molecules since the solute molecule is enveloped by solvent molecules only. Thermodynamic properties of solute molecules are highly dependent on the nature of solvent molecules. The variation of the thermodynamic properties with solvents will give us the characteristics of many interacting abilities of solute molecules, such as dispersion, induced dipole, dipole-dipole interaction, Coulomb force, charge transfer interaction, coordinate bond, and hydrogen bonding energies. Solution at infinite dilution is an ideal state for solute molecules. The properties of the solution are easily combined with current theories of solutions of molecules. Gas-liquid chromatography (GLC) provides the properties at infinite dilution on the basis of GLC conditions. Recently, GLC

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is a rapid and reliable method for determinination of accurate thermodynamic properties (1, 2, 4-12, 15-21).

In this paper, GLC was used here to study alkylbenzenes in three high-molecular-weight normal alkane solvents. Vaporliquid equilibrium constants at infinite dilution are obtained at five temperatures. Heat of solution from the gas phase to n-alkane solvent is also determined and can be given by a relation with structure of the solute molecule by a consistent set of parameters.

# **Experimental Section**

The gas chromatograph used was an Hitachi K53 equipped with a thermal conductivity detector. Column temperatures for measurement of retention volumes are 78.0, 84.0, 90.0, 96.0, and 102.0 °C. The column temperatures were maintained constant within 0.1 °C. Hydrogen was used as a carrier gas. The carrier gas flow rate was measured at room temperature with a soap-film flowmeter placed at the detector outlet. Inlet and outlet pressure of carrier gas was determined to 0.1 mmHg by mercury manometers.

Three stationary liquid phases were used: n-octacosane  $(n-C_{28}H_{58})$ , *n*-dotriacontane  $(n-C_{32}H_{66})$ , and *n*-hexatriacontane (n-C<sub>36</sub>H<sub>74</sub>). These phases were obtained from the Tokyo Chemical Industry Co., Ltd., Tokyo, Japan. n-C32H66 was purified three times by recrystallization from ethanol, while n-C<sub>28</sub>H<sub>58</sub> and n-C<sub>36</sub>H<sub>74</sub> were purified once similarly.

Benzene, toluene, ethylbenzene, p-, m-, and o-xylenes, cumene, and 1,3,5-, 1,2,4-, and 1,2,3-trimethylbenzenes were chosen for this study of alkylbenzenes and used without further purification.

The liquid for stationary phase was coated on to solid support material in the usual manner. The solid support used was Celite 545, 80-100 mesh, hexamethyldisilazane (HMDS) treated. The amount of each liquid phase on the support material was changed from 10 to 15 wt %. The coated support material was packed into 1 m X 3 mm i.d. stainless steel tubing. For reference, stationary-liquid-uncoated support material was also packed into different tubing. Each column was conditioned at 110 °C for about 24 h with a gentle flow of a carrier gas. The smallest detectable sample sizes were used at all times so as to ensure operation in the Henry's law region of solute con-

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